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Boundary-Layer Pressure Fluctuations at High Reynolds Numbers on a Free-Flight Test Vehicle

by

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BOUNDARY-LAYER PRESSURE FLUCTUATIONS AT HIGH REYNOLDS NUMBERS ON A FREE-FLIGHT TEST VEHICLE

by

D. R. Roberts

SUMMARY

Measurements have been made of the boundary-layer pressure fluctuations on the body of a free-flight aerodynamic test vehicle powered by a solid-fuel rocket motor. The vehicle reached a maximum Mach number of 2.2 with a maximum Reynolds number of about 215 millions.

Pressure spectra have been deduced, and have been found to compare reasonably with a theoretical spectrum for homogeneous isotropic turbulence.

The scale of the boundary-layer turbulence was found to fluctuate between 47% and 76% of the turbulence boundary-layer thickness over a range of Mach numbers from 1.5 to 2.2, while being essentially equal to 50% of this thickness over the range Ma = 2.0 to Ma = 2.2.

At Ma = 2.2 the root mean square boundary-layer pressure was equal to 0.0045 of the free stream dynamic pressure.

* Replaces RAE Technical Report 71033 - ARC 33171

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1 INTRODUCTION

A series of three free-flight aerodynamic test vehicles, designated the 'Shark'* series, was envisaged for the investigation of boundary-layer characteristics at high Reynolds number.

On Shark 1, measurements were made of skin friction drag by means of surface pitot tubes, and boundary-layer pressure fluctuations using a piezoelectric pressure transducer. On Shark 2, it was planned to make direct measurements of skin friction drag by means of a modified force-balance accelerometer, for comparison with the results obtained from Shark 1. Difficulties with the instrument resulted in the abandonment of the experiment. On Shark 3, measurements were made of boundary-layer pressure fluctuations employing two piezoelectric pressure transducers, one at a rear station similar to that used on Shark 1, and the other at a station positioned to give a Reynolds number half that at the rear station. This Report presents the pressure-fluctuation results from Shark 1.

2 DESCRIPTION OF TEST EXPERIMENT

2.1 Description of test vehicle

2.1.1 General

The vehicle consisted of a 4 calibre tangent circular ogival nose fitted to a cylindrical body 8 calibre long followed by a 2 calibre tail section, to the tail section were fitted three stabilising fins, radially equally distributed. The body and tail tube were constructed of resin bonded paper tube, and the nose was a glass fibre moulding. The vehicle was propelled by a non-separating solid-fuel rocket motor mounted inside the body. The overall length was about 233 inches (5.92 m) and the launch weight was about 1204 pour ds (546 kg). A general arrangement drawing is given in Fig.la.

2.1.2 Instrumentation

Transducers were installed to measure:

(a) Boundary-layer pressure fluctuations (special piezo-electric lead-zirconate bimorph transducer, ± 0.2 lbf/in² [± 1379 N/m²]).

*These vehicles were designed and constructed by Aerodynamics Department, RAE, and launched by them at the RAE Aberporth firing range.

(b) Pitot and static pressure differential at each of five surface pitot tubes (variable inductance pressure transducers).

(c) Acceleration normal to the vehicle body surface near the boundarylayer pressure measuring station (piezo-electric accelerometer).

(d) Accelerations normal and perpendicular to a given vehicle-body radius (variable inductance accelerometers).

The output signals of these transducers were translated through suitable sequence-switches and modulator units to frequency-modulate the transmission of three 465 MHz oscillator units. The signals from these oscillators were received and translated at the ground receiving station, and recorded on magnetic tape and oscilloscope record films.

One transmitter was used to telemeter the data from (b) and (d), to give measurements of vehicle accelerations and skin-friction. The second transmitter telemetered the output of (c) to give measurements of the environmental vibration local to the boundary-layer pressure-fluctuation transducer. The third transmitter telemetered the output from (a); the information derived from this transducer is considered in this paper.

2.2 Aerodynamic environment

The maximum Mach number achieved was about 2.2 $[q_{max} \simeq 6870 \ lbf/ft^2$ (330 kN/m²)].

The maximum local Reynolds number was about 215 millions. Maximum altitude was about 1456 ft (444 m).

2.3 Range facilities

In addition to the telemetry receiving and recording station, the firing range provided the following:-

(a) a number of observation posts suitably sited, equipped with highspeed cine cameras and kinetheodolites, which tracked the vehicle during flight, and afforded data from which was deduced velocity and altitude information.

(b) Radar cover to provide additional velocity and altitude data, particularly when the vehicle was beyond optical range of the observation posts, or obscured by cloud.

(c) A central source of timing pulses which were transmitted to all the data recording systems, and also initiated a firing sequence which resulted in the electrical pulse sent to the rocket motor igniter circuit. Measurement of the roll of the vehicle was also made by the telemetry receiving station.

The transmission of data from the boundary-layer pressure fluctuation transducer ceased after about 5 seconds of flight. This record deals with the data received during this time.

3 METHOD OF DATA ANALYSIS

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The signal representing the output of the boundary-layer pressure fluctuation transducer was frequency-translated after reception and recorded on $\frac{1}{2}$ inch magnetic tape at 60 inches per second. It was later played back at 7.5 inches per second, and re-recorded at 30 inches per second. Thus one second of vehicle flight time was represented by 240 inches of magnetic tape. This operation was carried out by Instrumentation and Ranges Department at the Royal Aircraft Establishment, Farnborough, using the equipment known as BRAMBLE¹. This analogue record was digitised by means of a high-speed digital recorder made to an RAE design by English-Electric/Leo/Marconi Ltd., the output from which was taken to a five-hole paper tape punch. For ease of handling, the record of the five seconds of flight time was divided into fifty punched paper tapes, each representing approximately 0.1 second of flight time, and designated runs 00 to 49, the designation representing the approximate time of the beginning of that section of data. Thus run 23, for instance, refers to the data received during the section beginning at about 2.3 seconds after the commencement of the flight.

Each of these fifty runs consisted of approximately 6170 discrete readings of pressure amplitude, expressed in units of analogue/digital converter output. For convenience in the computing process, these output units were retained throughout the analysis. The relationship between these units and standard physical units is given below:-

l pressure unit = 0.0031163 lbf/in² = 0.02149 kN/m² l(pressure)² unit = 9.7113 × 10⁻⁶ (lbf/in²)² = 0.4618 × 10⁻⁶ (kN/m²)².

The readings were spaced at intervals of 15.5 microseconds of flight time. Figs.2a, 2b and 2c show plots of pressure-readings for the first 5 milliseconds of each of the three runs, 32, 33 and 34.

Using the RAE Mathematics Department ICL 1907 computer, the data for each run was treated separately, the following operations being performed.

(a) The arithmetic mean of all the values in the run was found, and subtracted from each value in turn:-

 $y_t = (recorded value) - (arithmetic mean)$.

(b) The root mean square of the pressure fluctuations was evaluated using the expression

$$\sigma_{p} = \left[\frac{1}{(N+1)} \sum_{t=1}^{N+1} y_{t}^{2}\right]^{\frac{1}{2}}$$

where (N + 1) is the number of discrete values contained in the run.

(c) The autocorrelation function is defined as

$$R_{\tau} = \frac{1}{(N + 2 - \tau)} \sum_{t=1}^{N+2-\tau} y_t y_{(t+\tau-1)}$$

where $\tau = 1, 2, 3 \dots (M + 1)$

where M is the number of correlated intervals used in the spectral density analysis.

The expression was evaluated for $\tau = 1, 2, 3 \dots$ (100).

(d) The unsmoothed spectral density estimates were calculated from²:-

$$V_{k} = \Delta t \left\{ R_{1} + R_{(M+1)} \cos \left[(k-1)\pi \right] + 2 \sum_{\tau=2}^{M} R_{\tau} \cos \left[\frac{\left[(k-1)(\tau-1)\pi \right]}{M} \right] \right\}$$

where $k = 1, 2, 3, \dots$ (M + 1)

and ∆t is the time interval between each two consecutive discrete values.
Tabulated values of unsmoothed spectral density are presented in Table 1.
(e) The frequency to which this estimate refers is given by:-

$$f_k = \frac{(k-1)}{2\Delta tM}$$

Values of k with the corresponding values of frequency are presented in Table 2.

The expression in (d) and (e) were evaluated for $k = 1, 2, 3 \dots 100$.

(f) The estimates of hanning-smoothed and hamming-smoothed spectral densities are given by:-

$$w_{k} = 0.5 V_{1} + 0.5 V_{2} \qquad k = 1$$

= 0.25 (V_{k-1} + V_{k+1}) + 0.5 V_k 1 < k < M
= 0.5 V_M + 0.5 V_{M+1} k = M + 1

hamming smoothing

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$$Z_{k} = 0.54 V_{1} + 0.46 V_{2} \qquad k = 1$$

= 0.23 (V_{k-1} + V_{k+1}) + 0.54 V_k 1 < k < M
= 0.46 V_M + 0.54 V_{M+1} k = M + 1

These expressions were also evaluated for $k = 1, 2, 3 \dots 100$.

Autocorrelation functions plotted against τ , and smoothed and unsmoothed spectral density estimates plotted against k, are presented for a number of representative runs in Figs.6a to 6e. The maximum value of unsmoothed spectral density and the frequency f_m at which it occurred, were determined, for each run. It should be noted that f_m was taken as the frequency of the peak value of the unsmoothed spectral density, regardless of fluctuations about the mean path through the experimental points. In a few runs, two peaks of similar magnitudes occurred at fairly close frequencies, causing some doubt as to which represented the proper value of f_m . In such a case, reference was made to the plots of smoothed spectral density, which indicated one sensible peak value. The frequency of this value was taken as the value of f_m .

(g) For very high Reynolds numbers (non-dimensional viscosity function A = 0) it was shown³ that the maximum (pressure)² spectral density occurred at

where X, the non-dimensional frequency, was defined as

$$X = a\omega = \frac{2\pi fa}{U}$$

where ω was the wave-number, f the measured frequency, U the velocity of the stream relative to the measuring station, and a was a length to make wave-numbers non-dimensional.

The frequency at which the spectral density is a maximum is defined as ${\rm f}_{\rm m}$ so that

 $2.111 = \frac{2\pi f_m^a}{U}$

or

 $a = \frac{2.111U}{2\pi f_{m}}$.

The scale of the turbulence was defined as

and

$$\hat{x}_{u} = LUa$$

LU = 0.74677 when A = 0

If the frequency position of the maximum value of the experimental spectrum is made to coincide with that of the theoretical spectrum, the scale of the turbulence is therefore given by -

$$k_{\rm u} = \frac{2.111 \ (0.746770)}{2^{\tau} f_{\rm m}}$$
$$= \frac{0.250880}{f_{\rm m}} \ .$$

This expression was used to calculate the value of ℓ_{ij} for each run.

(h) The local Reynolds number R_x was calculated for each run, assuming that the effective start of the turbulence was at a distance x upstream of the measuring station, equal to the axial distance of the station from the nose of the test vehicle, and using the data of velocity and altitude.

(1) Taylor⁴ has concluded that a reasonable estimate of the turbulence thickness $\$_T$ is equal to about 80% of the total boundary-layer thickness \$ predicted by Spalding⁵, so that

$$R_{s_{T}} = 0.8 R_{s}$$

The relationship between R_{s} and R_{x} has also been defined by Spalding, and using this relationship together with the expression

$$F_{T} = \left(\frac{0.8 R_{\$}}{R_{x}}\right)x$$

the values of $\ensuremath{\$_{\tau}}$ were calculated for each run.

*

(j) The ratio $\ell_{\rm u}/\$_{\rm T}$ was then calculated for each run.

3.1 For runs 00 to 20, (Ma = 0 to Ma ~ 1.2) the signal/noise ratio of the recorded data was too small for the effective determination of a peak in the plot of spectral density against frequency factor. Under these conditions, a spurious peak imposed by the reading frequency was often apparent, at the lowest frequency used. This would lead, of course, to a false value of ℓ_u being obtained if this frequency was used in the calculation of the scale of the turbulence in the expression given in 3(g) above. For this reason, the values of the computed spectral density estimates for these runs have not been tabulated, or used in analysis.

The data for runs 21 to 24 inclusive have not been fully analysed because the signal/noise ratios, though higher than those for earlier runs, were still too low for really adequate identification of a significant peak. The data have been included in Table 1 merely for possible use in transonic studies.

The spectral density estimates for runs 21 to 49 are tabulated for a range of k = 1 to k = 100, in Table 1, and for runs 25 to 49 (Ma ≈ 1.5 to Ma ≈ 2.2) the values of the scale of the turbulence ℓ_u and of the ratio $\ell_u/\$_T$ were calculated. For runs 33 and 49, however, the spectral density plots show very low mean values and low signal/noise ratios, and in these cases the results have not been tabulated. It is suspected that these two runs were affected by transient telemetry faults, and immediately after run 49, the transmission of data effectively ceased.

The values of ℓ_u obtained are plotted against time in Fig.7a and against Mach number in Fig.7b. The positions of runs 33 and 49 are indicated, but no values plotted. Similarly, positions are indicated, but no values plotted, for runs 33 and 49 in the plots of $\ell_u/\$_T$ against time (Fig.10) and against Mach number (Fig.11).

Values of all the environmental experimental data and derived quantities are presented in Table 3, and the vehicle trajectory data in Figs.3, 4 and 5. 3.2 To reinforce the conclusions derived from the analyses described above, the computing technique was modified to accept the pressure readings for five consecutive runs and use them as the data for one run.

The readings for runs 41 to 45 were used as the data for this collective run, which thus contained about 30800 discrete readings, separated by intervals $\Delta t = 15.55$ microseconds. The mean velocity and mean altitude over the

considered time span were used to derive Reynolds numbers and hence ${}_{T}$, and the autocorrelation function, unsmoothed and smoothed spectral density estimates, scale ${}_{u}$, ratio ${}_{u}/{}_{T}$ and root mean square pressure σ_{p} were obtained as for the individual runs.

The unsmoothed spectral density estimates are shown in graphical form in Figs.9 and 12c, and the experimental environmental data are given in Table 4.

4 DISCUSSION OF RESULTS

The pressure-transducer used to measure the wall pressure was carefully mounted and inspected to ensure that the measuring diaphragm was flush with the exterior skin of the vehicle. The diaphragm diameter was 0.125 inch (3.175 mm) and the maximum projection above the surrounding surface was 0.003 inch (0.076 mm).

It was arbitrarily assumed that pressure-fluctuations having a wavelength of less than ten times the diameter of the transducer diaphragm would be attenuated to such an extent that they should be regarded as being unreliable, i.e. the highest usable frequency of pressure-fluctuation is given by

$$f = \frac{v}{10d}$$

where v is the vehicle velocity and d is the diaphragm diameter. For the range of vehicle velocities considered (i.e. 1692 ft/s to 2444 ft/s) and the transducer diaphragm diameter of 0.01042 ft, the value of f lies in the range 16.24 kHz to 23.45 kHz, corresponding to values of k = 51 to k = 73. Since all the significant peaks in the plots of spectral density against k occur well below k = 40, it is assumed that errors due to spatial resolution of the transducer may be ignored.

The data obtained from the piezo-electric accelerometer (c) showed that although the vibration local to the piezo-electric pressure transducer was of considerable amplitude in the frequency-range 100 to 300 Hz, the amplitude at higher frequencies was very much smaller, and would have negligible effect on the accuracy of the pressure data.

4.1 <u>Comparison of experimentally determined spectra with a theoretical</u> <u>spectrum for homogeneous isotropic turbulence</u>

In Ref.3, Taylor tabulated a theoretical $(pressure)^2$ spectrum for very high Reynolds numbers, (A = 0) in a non-dimensional form, and this was used as a basis for comparison with the experimental spectra. Figs.12a, b and c show respectively the experimental spectra determined for run 35 (Ma = 2.060), run 45 (Ma = 2.188) and the collective run described in section 3.2 (mean Ma = 2.193). The strong resemblance between these three spectra confirms a reasonably consistent state of turbulence throughout this portion of the available data.

The theoretical spectrum has been superimposed on each of these three figures. The full line indicates the spectrum soplaced as to have its maximum value at the same frequency as that of the maximum experimental value, and to represent the same value of $(\sigma_p)^2$ as the experimental data. The broken line indicates the spectrum placed so as to obtain a good fit with the experimental points over the high frequency range, whilst retaining the same value of $(\sigma_p)^2$.

Although the test vehicle was subject to longitudinal acceleration ranging, over the time span considered, from +22 g to -4 g, and could not therefore be regarded as being in a steady state and thus not directly comparable with the theoretical case, nevertheless a surprisingly good agreement was apparent. It may be noted that the collective run considered in Fig.12c was derived from data obtained during the time span centred on the 'zero g' point of the flight.

4.2 The scale of the turbulence, ℓ_{u}

The values of ℓ_u given in Tables 3 and 4 are based on the values of f_m indicated by the coincident peaks of the experimental points, and the full-line theoretical spectrum. If the values of ℓ_u were to be based on values of f_m given by the peak of the broken-line spectrum, they would be greater by a factor of about 2.6. The values of ℓ_u obtained from the two values of f_m for the three runs illustrated in Figs.12a, b and c are given below:

Run	(Full-line) ^l u(f)	(Broken-line) ^L u(b)	$\frac{{}^{\ell}u(b)}{{}^{\ell}u(f)}$
35	0.09849 ft (0.03002 m)	0.2616 ft (0.0797 m)	2.66
45	0.08953 ft (0.02729 m)	0.2266 ft (0.0691 m)	2.53
Collective	0.08970 ft (0.02734 m)	0.2266 ft (0.0691 m)	2.53

The true value of ℓ_u associated with the pressure fluctuations due entirely to turbulence probably lies somewhere between these two extremes.

The values of ℓ_u presented throughout this paper are obtained from calculations based on the frequency of the experimental peak values. However, since the flight of the test vehicle was of such short duration, it is possible that the turbulence was not fully established and that the values of ℓ_u given are somewhat low.

The plot of ℓ_u against elapsed time Fig.7a indicates a mean value of about 0.09 ft (27.4 mm), the individual points showing very little deviation in level after 3.3 seconds. Fig.7b, a plot of ℓ_u against Mach number, again shows very little deviation from the mean above a value of Ma = 2. As the value of $\$_T$ is sensibly constant over the considered range of Mach number, the values of $\ell_u/\$_T$ shown in Figs.10 and 11 also exhibit this constancy.

The mean value of ℓ_u for the runs $41 \rightarrow 45$ is 0.09245 ft (28.2 mm) and ℓ_u for the collective run is 0.0897 ft (27.3 mm). The mean value of $\ell_u/\$_T$ for the runs $41 \rightarrow 45$ is 0.5275, and the value for the collective run is 0.5128. Therefore, the results obtained from the collective run may be assumed to be representative of those from the individual runs.

Taylor⁴ cites an example of experimental results derived from an investigation of turbulence at subsonic velocities, which are compared below with the results from the present investigation.

	Ma	^R \$ _T	\$ _T	\$	^l u	$\frac{u}{*}$ T	σ <u>p</u> q
Taylor	0.176	3.8×10^4	0.328 ft (100 mm)	0.410 ft (125 mm)	0.0975 ft (29.7 mm)	0.297	0.0056
Shark	2.2	2 × 10 ⁶	0.175 ft (53.3 mm)	0.219 ft (66.8 mm)	0.0897 ft (27.3 mm)	0.513	0.0045

It is seen that although the Mach number has increased from the first case to the second by a factor of 12.5, and the turbulence Reynolds number by a factor of 53, the scale of the turbulence has remained almost the same, and the ratio $\ell_{\rm u}/\$_{\rm T}$ has increased only by a factor of 1.73.

5 CONCLUSIONS

Measurements of boundary-layer pressure fluctuations have been made for a short flight-time on a free-flight aerodynamic test vehicle at Mach numbers in the range 1.5 to 2.2, and local Reynolds numbers of 1.51 to 2.15×10^8 . Autocorrelation functions and (pressure)² spectral densities have been determined and are presented in tabular and graphical form. Representative experimental spectra are compared with a theoretical spectrum for homogeneous isotropic turbulence and are found to be in reasonable agreement.

Values have been obtained of the scale of the turbulence, and the ratio of scale of the turbulence to turbulence boundary-layer thickness, and these are similarly presented. Over the considered range of Ma = 1.5 to Ma = 2.2, the scale of the turbulence was found to fluctuate between 0.47 and 0.76 of the turbulence boundary-layer thickness. At Ma = 2.2 the scale of the turbulence was about 0.51 of the turbulence boundary-layer thickness, and the root mean square pressure was 0.0045 of the free-stream dynamic pressure.

Table.1 Page 1 of 8 pages

DENS	ITY OF (JENCY FA	PRESSUR	E SQUAR ". (100	ED X 10 PRESSU	000) SP RE UNIT	ECTRUM S = $0.$	WITH RES 3116 LB/S	SPECT TO SQ INCH)
ĸ			R	UN NUME	ER		· · · · · · · · · · · · · · · · · · ·	
	21	22	23	24	25	26	27	
123456789012345678901234567890123456789012334567890123901234567890123456789012345678901234567890123456789012345678901234567890123456789012345678901234567890123456789012345678901234567890012345678900123456789001234567890012345	4990291425227539825551504971624627904777394155590739991	34539597500216227261777588280121374619053628966675237	4343544434343422232233221212121111111111	6207464720193388343437183100130041763565134084933541	26951778336454464444333212222121212212111111118995938314 9511783364544930862237712221212121211111111899593838314 1888867	23222221112212111111111111111111111111	25204762670772599333680015036697112111111111111111111111111111111111	

DENS FREQI	ITY OF (JENCY FA	PRESSUR CTOR [*] k	E SQUAR *. (100	ED X 10 PRESSU	000) SP RE UNIT	ECTRUM S = 0.3	WITH RES 116 LB/S	PECT TO Q INCH)
k			R	UN NUMB	ER			
	21	22	23	24	25	26	27	
1234567890123456789012345678901234567890123456789012345678901234567890 1012345678901234567890123456789012345678901234567890	3890922392185544874775748533032484869270255072897765	41359137230587434909965329381690591921637948568800	44 53552444432322212122221212121212121212121212	4733288865031883104307666613670881737371586136570817	7554272322790509111161586856900339770744976869274780	86755848755443333533324312232222321574221211221 1574221211221 12211221	980493241851981663192616386040682711742777393748326 16777779765547775545446386040682711742777393748322	

Table 1. Page 2 of 8 pages

Table.1. Page 3 of 8 pages

DENS FREQL	DENSITY OF (PRESSURE SQUARED X 10000) SPECTRUM WITH RESPECT TO FREQUENCY FACTOR "K". (100 PRESSURE UNITS = 0.3116 LB/SQ INCH)											
k	RUN NUMBER											
	28	29	30	31	32	33	34					
123456789012345678901234567890123456789012334567890123444444444444444444444444444444444444	$\begin{array}{c} 397\\ 1154\\ 1003\\ 9870\\ 98220\\ 12204\\ 10264\\ 9870\\ 122064\\ 10265\\ 10266\\ 10265\\ 10266\\ $	$\begin{array}{c} 101\\ 16929\\ 1765426\\ 56478\\ 162216\\ 111\\ 1216226\\ 12333\\ 1216226\\ 1212223\\ 1212223\\ 1212223\\ 1212223\\ 1212223\\ 1212223\\ 1212223\\ 1212223\\ 1212223\\ 1212223\\ 121$	$\begin{array}{l}9592\\ 95762\\ 21222116648826807625962772666555764449778219913957720001879558769212122122126666557664497821395773620018795587958795877652211221221666655576444997821391395773620018799587955876666555766449978211913957736200187995879558766666555766467622169879958795587666665557664676221698799572000187995720001879957200018799587955876666655576646762216987995879573665001879957766666555766467622169879957200018799572000187995720001879957200018799587957666665557664676221698799572000187995720001879957200018795576666555766467622169879957200018799572000187995795766666555766467622169879957200018799572000187995720001879957200018799579576666655576646762216987995879576666655576676666655576666655576676666555766766$	$\begin{array}{c} 95369\\ 9537740\\ 95322288644\\ 193517469\\ 323922288644\\ 193517469\\ 323922222\\ 187732555754444544\\ 153322322222222222222222222222222222222$	$\begin{array}{c} 1299149792323232323232323232323232323232323232$	4465433944243232321211111111111 981038821698852217774582587943400113771126506374987843434	$\begin{array}{c} 436\\ 11280857769710322122231444522267119847655555556112222231244567132222671129847675555555611522333332211122222312445561119864119087866694495561150380723212122223124455221112222231244552211122222312445522111222223124455221112222231244552211122222312445522111222223124455222231244552222312422267112222231222231244552222312222312222231222223122222312222231222223122222312222231222222$					

DENS I FREQU	TY OF (ENCY FA	PRESSUR CTOR "k	E SQUAR . (100	ED X 10 PRESSU	000) SPI RE UNIT	ECTRUM S = 0.3	WITH RE 116 LB/	SPECT TO SQ INCH)
 ĸ		· · · · · · · · · · · · · · · · · · ·	R	UN NUMB	ER			
	28	29	30	31	32	33	34	
55555555556666666666666777777777788888888	121 111 128 78 96 87 55 65 67 63 55 440 53 76 34 82 87 940 35 1630 890 48 8993 114 33 232 222 312 121 23 23 222 312 121 23 23 23 222 312 121 23 23 23 23 23 23 23 23 23 23 23 23 23	124523092648210343885747018424516009532484103602127	14043320073468068670774397022624584473313879120574433	21111111111111111111111111111111111111	1904122112121111111111111111111111111111	333323232122111111111111111111111111111	2212212212211111112111211111 1122212221	

Table 1. Page 4 of 8 pages

Table	1.	Page	5 of	8	pages
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DENS FREQU	ITY OF JENCY FA	(PRESSUR ACTOR #1	RE SQUAN K ^N . (10)	RED X 1 D PRESS	0000) SI URE UNI	PECTRUM TS = 0.	WITH RE 3116 LB/	SPECT TO SQ INCH)
ĸ		······································]	RUN NUM	BER			
	35	36	37	38	39	40	41	
123456789012345678901234567890123456789012334567890123444444444444444444444444444444444444	$\begin{array}{c} 664\\ 941\\ 195020\\ 5020\\$	887233384384444502818115558655298153438079053949534 8872335843844459750281817558629815334380790539499534	$\begin{array}{c} 6876\\7821\\997399944936800\\8272217508272334323211537968673919748665466663298\\429756800851123729057686730494397919748665466663298\\429759867309494397919748665466663298\\42975986730494397919748665466663298\\42975986730494397919748665466663298\\42975986730494397919748665466663298\\429759867304943979197486655466663298\\429759867304943979197486655466663298\\429759867304943979197486655466663298\\429759867304943979197486655466663298\\429759867304943979197486655466663298\\4297598673049439791974866554666653298\\4297598673049439791974866554666653298\\42975986730649439791974866554666653298\\42975986730649439791974866554666653298\\429759867306494397991974866554666653298\\429759867306494397991974866554666653298\\429759867306494685966554666653298\\429759667596675966759667596665546666554666655986655666655666656566565$	$\begin{array}{c} 9019759820095963143749005566781724999014373925630498356132\\ 90197582009596314374900556678172499654959654455724508356132\\ 9019758200959631437490055667817249990143739256540498356132\\ 90198420095963143749005566781724999014373925630498356132\\ 901984200959631437490055667817249990143739256540498356132\\ 90198420095967847490055678172499901437392567345489900182\\ 90198420095967847490055678172499901437392567345489900182\\ 90198420095967847392567817249990143739256784598356132\\ 901984200959678473925678172499901437392567845724993332\\ 9019842009596784500498356132\\ 90198420095967847392567849990143739256738056784993333\\ 90198420095967847392567849990143739256738056784993333\\ 901984200959678459049837926678678678678678678678678678678678678678$	845299341111111211122233333322219188969755865748285732477759 111111221118881034831615002564085732477759	$\begin{array}{c} 164516985110184665310759122249269506732811967454544949786873\\ 111112224432214926950673281119678765737545454493286873\\ 11112224392512266950673281119678765737545454493286873\\ 111122244322812266950673281119678765737545454493286873\\ 111122244322812266950673281119678765737545454493286873\\ 1111222443228122669506732881199671226497868873\\ 111122244322812669506732881199671226497888873\\ 11112224432281266950673288119967122649788765732881199671226392811\\ 11112224432281266950673288119967122649788765732881199671226697832811196788766873288119967122669783281119678876573754545492668732881196787876573288119967122649288873\\ 1111267878786887388888738888887388888873888888873888888$	$\begin{array}{c} 1103\\ 989447862424411122232379333122111159334794885536346553533479488555363465535334711111111112223232333122214111292333479488555362433465536347424441414141412223232333342221441414141414141$	

Table 1. Page 6 of 8 pages

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DENS I FREQU	DENSITY OF (PRESSURE SQUARED X 10000) SPECTRUM WITH RESPECT TO FREQUENCY FACTOR "K". (100 PRESSURE UNITS = 0.3116 LB/SQ INCH)										
ĸ			F	RUN NUME	BER	· · · · · · · · · · · · · · · · · · ·					
	42	43	44	45	46	47	48	49			
12345678901234567890123456789012345678901233456789012345678901234567890	$\begin{array}{c} 14575\\ 14982\\ 0505\\ 050$	$\begin{array}{c} 10118744526801773165480732668754312999879996046291300992 \end{array}$	$\begin{array}{c} 1011111111111111111111111111111111111$	$\begin{array}{c}9122315891845004100873897449890272882460103055648888199941\\122552439116439164648881999902724527767444609388199941\\\end{array}$	$\begin{array}{c} 426959811585399445887066264084413117766675666344013035499462294 \\ 111111111111111111223332111111776676005273611303549465738995 \\ 1111111111111111111111111111111111$	264731122977867514058735642511000534337869344384544370	$\begin{array}{c} 41529\\ 7619222437104344359319422076201686853540953524981485642\\ 1211212111212111112121111211211111111$	$\begin{array}{c} 5252\\ 1928\\ 1371928\\ 1371946\\ 185525\\ 196522\\ 1121216\\ 1121248\\ 1966525\\ 1855258\\ 165226\\ 1121248\\ 196699\\ 1229994\\ 067922\\ 999933\\ 999986\\ 637524\\ 037524\\ 03866\\ 1057752\\ 630\\ 1057752\\ 0395\\ 1057752\\ 0395\\ 1057752\\ 0395\\ 03866\\ 03752\\ 0366\\ 0366\\ 0366\\ 03752\\ 0366\\ $			

Table.1. Page 7 of 8 pages

Table.1. Page 8 of 8 pages

Table.2 Frequency and corresponding frequency factor 'k'.

k FREQUENCY H		k	FREQUENCY (Hz)	k	FREQUENCY (Hz)	ĸ	FREQUENCY (Hg)
1	0.00	26	8119.77	51	16239.55	76	24359.32
2	324•79	27	8444.57	52	16564.34	77	24684.12
3	649.58	28	8769.36	53	16889.13	78	25008.91
4	974•37	29	9094.15	54	17213.92	79	25333.70
5	1299.16	30	9418.94	55	17538.71	80	25658.49
6	1623.95	31	9743•73	56	17863.50	81	25983.28
7	19 48•75	32	10068.52	57	18188.30	82	26308.07
8	2273.54	33	10393.31	58	18513.09	83	26632.86
9	2598.33	34	10718.10	59	18837.88	84	26957.65
10	2923.12	35	11042.89	60	19162.67	85	27282.44
11	3247•91	36	11367.68	61	19487.46	86	27607.23
12	3 572 .7 0	37	11692.48	62	19812.25	87	27932.03
13	3897•49	38	12017.27	63	20137.04	88	28256.82
14	4222.28	39	12342.06	64	20461.83	89	28581.61
15	4547.07	40	12666.85	65	20786.62	90	28906•40
16	4871.86	41	12991.64	66	21111.41	91	29231.19
17	5196.66	42	13316.43	67	21436.21	92	29555 .9 8
18	5521.45	43	13641.22	68	21761.00	93	29880.77
19	5846.24	44	13966.01	69	22085.79	94	30205.56
20	6171.03	45	14290.80	70	22410.58	95	30530.35
21	6495.82	46	14615.59	71	22735•37	96	30855.14
22	6820.61	47	14940•39	72	23060.16	97	31179.94
23	7145.40	48	15265.18	73	23384.95	98	31504.73
24	7470.19	49	15589.97	74	23709.74	99	31829.52
25	7794•98	50	15914.76	75	24034.53	100	32154.31

			t																						
48	47	46	45	44	43	42	41	40	39	36	37	36	35	34	33	32	31	30	29	28	27	26	25		RUN
2409	2418	2426	2434	2440	2444	2444	2439	2429	2413	2392	2365	2333	2295	2241	2202	2149	2091	2030	1965	1899	1830	1762	1692	(ft/s)	VELOC ITY
1412	1368	1324	1280	1236	1192	1148	1104	1060	1018	976	944	868	850	800	766	726	682	640	606	572	538	504	470	(ft)	HE I GHT
2.1668	2.1746	2.1814	.2.1883	2.1933	2.1966	2.1962	2.1914	2.1821	2.1674	2.1482	2.1237	2.0944	2.0601	2.0113	1.9761	1.9282	1.8759	1.8209	1.7624	1.7030	1.6409	1.5797	1.5168		MACH
010_775	211.787	212.713	213.642	214.396	214.976	215.206	214.995	214.342	213.148	211.509	209.284	206.836	203-556	199.010	195.709	191.185	186.226	180.980	175-331	169.584	163.559	157.613	151.478	-6 (x10 ·)	×R
2.612	2.623	2.634	2.644	2.653	2.659	2.662	2.660	2.652	2.639	2.620	2.595	2.568	2.531	2.479	2.442	2.390	2.334	2.274	2.209	2.143	2.074	2.005	1.934	-6 (X10 ⁻)	- ²⁷

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REFER TO SECTION 3.1

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0.181 0.176 0.175 0.175 0.175 0.175 0.175 0.176 0.176 0.176 0.176 0.177 0.178 0.178 0.178 0.179 0.180 0.180 0.176 0.175 0.175 0.175 0.175 0.175 0.177 (ft) ୶ଡ଼ а ^ж 21 22 20 21 20 22 19 19 19 44 $\mathbf{1}_{\mathbf{3}}$ 14 17 14 21 20 21 22 21 27 20 3 3 * ٠ 0.41451 0.42521 0.09150 0.12635 S.D. MAX 0.39302 0.37992 0.41882 0.36527 0.32889 0.39923 0.43791 0.39385 0.79481 0.56015 0.50926 0.57465 0.40512 0.40361 0.20259 0.06800 0.06335 0.31045 0.34427 ٠ ٠ 0.10012 0.09424 0.09439 0.09319 0.09725 0.09849 0.09222 0.13067 0.11676 0.12223 0.10873 0.08506 0.10054 0.09370 0.08953 0.09439 0.08971 0.09875 0.08699 0.09111 0.13459 0.09303 0.09830 SCALE (ft) * ۰

Table 3

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*

SCALE/\$	R.M.S. PRESSURE	RUN
	(Lbf/in)	
0.5559	0.0964	25
0.4720	0.1015	26
0.6056	0.1148	27
0.6825	0.1332	28
0.6541	0.1544	29
0.7342	0.1842	30
0.7578	0.2152	31
0.5207	0.2355	32
*	0.0835	33
0.5163	0.2184	34 34
0•5592	0.2232	Э 5
0.5693	0.2269	36
0.4952	0.2171	37
0.5541	0.2112	38
0-5314	0.2094	39
0.5634	0.2126	40
0.5119	0.2100	41
0.5387	0.2261	42
0.5386	0.2130	43
0.5376	0.2167	44
0.5106	0.2111	45
0.5342	0.2146	46
0.5602	0.2321	47
0.5307	0.2475	48
*	0.2229	49

	IMPER IAL	UN ITS	s.1. UI	NITS
MEAN VELOCITY	2440	ft/s	743.7	m/s
MEAN HEIGHT	1192	ft	363•3	E
MEAN MACH NUMBER	2.193			
r (x10 ⁻⁶) x	214.64			
R (X10 -) \$ T	2•656			
0 ₩	0.175	ft	0•053	٤
×Ε	22			
S.D. MAX	0•3562			
SCALE	0.0897	ft	0.027	E
SCALE/\$ T	0.5128			
R.M.S. PRESSURE	0.2153	2 Lbf/in	1.4844	kN/m

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Table.4 Environmental data and derived quantities for collective run.

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A	non-dimensional viscosity function
LU	l _u /a
Ma	Mach number
М	number of correlated intervals used
N	number of intervals in one run
R _{\$}	boundary-layer Reynolds number
^R \$ _T	turbulence Reynolds number
R _x	local Reynolds number at distance x from the effective start of turbulence
R _τ	autocorrelation function for interval τ
U	stream velocity relative to measuring station
$\left. \begin{smallmatrix} v_k \\ s D_k \end{smallmatrix} \right\}$	unsmoothed spectral density estimate at frequency f_k
SD max	maximum value of unsmoothed spectral density estimate for one run
W.k	hanning-smoothed spectral density estimate at frequency f_k
Х	ωα
Z _k	hamming-smoothed spectral density estimate at frequency f_k
а	a length to make wave numbers non-dimensional
fk	frequency corresponding to frequency factor k (see Table 2)
f _m	frequency at which the maximum spectral density occurs
k	integer included between 1 and (M + 1); the frequency factor corresponding to frequency f_k
k _m	frequency factor corresponding to f m
^l u	the scale of the turbulence
∆ _t	time interval separating each discrete value y_t from the next
У _t	recorded discrete pressure value at time t, referred to the mean value for that run
\$	total boundary-layer thickness
\$ _Т	turbulence thickness

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SYMBOLS (Contd)

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τ	integer included between 1 and $(M + 1)$
ω	wave number = 2π /wavelength
σp	root mean square pressure
q	free stream dynamic pressure

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Fig. la General arrangement of `Shark l'

Fig1(b) Side view photograph of 'Shark 1'

Fig.2b Plot of pressure against time for the first 5 milliseconds of run 33

Fig.2c Plot of pressure against time for the first 5 milliseconds of run 34

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Fig. 7a Scale/elapsed time

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Fig 7b Scale/Mach number

Fig.8 b \$_T / Mach number

Fig 9 Collective run-unsmoothed spectral density/k

Fig. 12 a Plot of unsmoothed spectral density (USSD) against frequency, for run 35

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Fig 12b Plot of unsmoothed spectral density (USSD) against frequency, for run 45

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DETACHABLE ABSTRACT CARD

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1	ARC CP No 1208 March 1971 Roberts D. P.		532,526,4 533 6 048 2 533 665 5	ARC CP No 1208 March 1971 Roberts, D. R	532,526,4 533.6 048 2 533 665 5
1	BOUNDARY-LAYER PRESSUR AT HIGH REYNOLDS NUMBE FLIGHT TEST VEHICLE	E FLUCTUATIONS RS ON A FREE-	533 6.011 12 533 6 011 5	BOUNDARY-LAYER PRESSURE FLUCTUATIONS AT HIGH REYNOLDS NUMBERS ON A FREE- FLIGHT TEST VEHICLE	533 6 011 12 533 6 011 5
 	Measurements have been made of a free-flight aerodynamic test veh reached a maximum Mach numbe 215 millions	the boundary-layer pressure the boundary-layer pressure to the powered by a solid-fuel ror of 2 2 with a maximum Rey	fluctuations on the body of ocket motor The vehicle molds number of about	Measurements have been made of the boundary-layer pressul a free-flight aerodynamic test vehicle powered by a solid-fue reached a maximum Mach number of 2 2 with a maximum 1 215 millions	re fluctuations on the body of I rocket motor The vehicle Reynolds number of about
1	Pressure spectra have been deduce theoretical spectrum for homogen	ed, and have been found to co neous isotropic turbulence	empare reasonably with a	Pressure spectra have been deduced, and have been found to theoretical spectrum for homogeneous isotropic turbulence	compare reasonably with a
; ; ;	The scale of the boundary-layer t of the turbulence boundary-layer while being essentially equal to 50	urbulence was found to fluctu thickness over a range of Mac 0% of this thickness over the 1	tate between 47% and 76% th numbers from 1.5 to 2 2, range Ma = 2.0 to Ma = 2 2	The scale of the boundary-layer turbulence was found to flu of the turbulence boundary-layer thickness over a range of M while being essentially equal to 50% of this thickness over th	ctuate between 47% and 76% fach numbers from 1.5 to 2 2, he range Ma = 2.0 to Ma = 2.2
1	At Ma = 2.2 the root mean squa free stream dynamic pressure	re boundary-layer pressure wa	as equal to 0 0045 of the	At Ma = 2 2 the root mean square boundary-layer pressure free stream dynamic pressure	was equal to 0 0045 of the
	532 526 4 file tange Ma = 2 0 to Ma = 2 2 file tange Ma = 2 0 to Ma = 2 0 to Ma = 2 2 file tange Ma = 2 0 to Ma	SSURE FLUCTUATIONS JIMBERS ON A FREE. Ide of the boundary-layer pre- st vehicle powered by a solid- immber of 2 2 with a maximu mogeneous isotropic turbulen mogeneous isotropic turbulen ayer thickness over a tange over li to 50% of this thickness over aguate boundary-layer pressi reserved	ARC CP No 1208 March 1971 Roberts, D R ATCHT TEST VEHICLE FLIGHT TEST VEHICLE Pressure spectra have been ma theoretical spectrum for hor theoretical spectrum for hor theoretical spectrum for hor treached a maximum Mach m Pressure spectra have been ma theoretical spectrum for hor theoretical spectrum for		
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